

$$1. \quad H = \frac{L_x^2 + L_y^2}{2I_1} + \frac{L_z^2}{2I_2}, \quad I_1 \gg I_2$$

$$= \frac{L_x^2 + L_y^2 + L_z^2}{2I_1} + L_z^2 \left(\frac{1}{2I_2} - \frac{1}{2I_1} \right)$$

$$H = \frac{L^2}{2I_1} + L_z^2 \left(\frac{1}{2I_2} - \frac{1}{2I_1} \right)$$

This has eigenvalues $\frac{l(l+1)\hbar^2}{2I_1} + m^2\hbar^2 \left(\frac{1}{2I_2} - \frac{1}{2I_1} \right)$

or $\frac{[l(l+1) - m^2]\hbar^2}{2I_1} + \frac{m^2\hbar^2}{2I_2}$

In limit $I_1 \gg I_2$, this is approx. $\frac{m^2\hbar^2}{2I_2}$.

2.

$$\text{Suppose } \langle l_x^2 \rangle = \langle l_y^2 \rangle = 0$$

$$\text{Then, } \langle L^2 \rangle = \langle L_z^2 \rangle$$

If $|l, m\rangle$ is an eigenstate of L^2, L_z , this implies

$$\langle l, m | L^2 | l, m \rangle = m^2 \hbar^2 \quad \text{or} \quad \langle l, m | L_z^2 | l, m \rangle = m^2 \hbar^2$$

This can only be true if $l = 0$.

3.

$$Y_{00} = \frac{1}{\sqrt{4\pi}} \quad (\text{constant})$$

$$Y_{10}(\theta, \varphi) = \sqrt{\frac{3}{4\pi}} \cos\theta = \sqrt{\frac{3}{4\pi}} \frac{z}{r} \quad z = r \cos\theta \quad (\text{geometry})$$

$$Y_{1\pm 1}(\theta, \varphi) = \mp \sqrt{\frac{3}{8\pi}} \sin\theta e^{\pm i\varphi}$$

$$= \mp \sqrt{\frac{3}{8\pi}} \sin\theta (\cos\varphi \pm i \sin\varphi)$$

$$= \mp \sqrt{\frac{3}{8\pi}} (\sin\theta \cos\varphi \pm i \sin\theta \sin\varphi)$$

$$= \mp \sqrt{\frac{3}{8\pi}} (x \pm iy)$$

$$Y_{20} = \sqrt{\frac{5}{16\pi}} (3\cos^2\theta - 1) = \sqrt{\frac{5}{16\pi}} \frac{3z^2 - r^2}{r^2}$$

$$Y_{2\pm 1} = \mp \sqrt{\frac{15}{8\pi}} (\sin\theta \cos\theta) e^{\pm i\varphi}$$

$$= \mp \sqrt{\frac{15}{8\pi}} \cos\theta (\sin\theta \cos\varphi \pm i \sin\theta \sin\varphi)$$

$$= \mp \sqrt{\frac{15}{8\pi}} \frac{z(x \pm iy)}{r^2}$$

$$Y_{2\pm 2} = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2\theta e^{\pm 2i\varphi}$$

$$= \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2\theta (\cos 2\varphi \pm i \sin 2\varphi)$$

$$= \sqrt{\frac{15}{32\pi}} \sin^2\theta (\cos^2\varphi - \sin^2\varphi \pm 2i \sin\varphi \cos\varphi)$$

$$= \sqrt{\frac{15}{32\pi}} \left[(\sin\theta \cos\varphi)^2 - (\sin\theta \sin\varphi)^2 \pm 2i (\sin\theta \sin\varphi)(\sin\theta \cos\varphi) \right]$$

$$= \sqrt{\frac{15}{32\pi}} \frac{(x \pm iy)^2}{r^2}$$

4.

$$\psi(x, y, z) = C(xy + yz + zx) e^{-\alpha r^2}$$

From problem 3,

$$xz = r^2 \sqrt{\frac{2\pi}{15}} (Y_{2-1} - Y_{21})$$

$$yz = i r^2 \sqrt{\frac{2\pi}{15}} (Y_{21} + Y_{2-1})$$

$$xy = -i r^2 \sqrt{\frac{2\pi}{15}} (Y_{2-2} - Y_{22})$$

$$\psi(r, \theta, \phi) = C \sqrt{\frac{2\pi}{15}} \left\{ (1+i) Y_{2-1} - (1-i) Y_{21} - i(Y_{22} + Y_{2-2}) \right\} r^2 e^{-\alpha r^2}$$

Since all Y_{lm} of form Y_{2m} , if a measurement of L^2 is made, the only possible eigenvalue is $\hbar^2(2+1)k^2 = 6\hbar^2$. $P(l=2) = 1$

If L_z is measured, possible eigenvalues are $\hbar k, \pm 2\hbar$.

$$\left. \begin{aligned} P(m=-2) &= P(m=2) = 1/6 \\ P(m=-1) &= P(m=1) = 1/3 \end{aligned} \right\} \quad 1/6 + 1/6 + 1/3 + 1/3 = 1$$

5.

$$V = \begin{cases} 0 & r > a \\ \infty & r < a \end{cases}$$

$$\begin{aligned} \langle \vec{r} | \psi^{(+)} \rangle &= e^{i\vec{k} \cdot \vec{r}} + \frac{e^{ikr}}{r} f(\vec{k}, \vec{k}') \\ &= \frac{1}{(2\pi)^{3/2}} \sum_{\ell} i^{\ell} (2\ell+1) A_{\ell}(r) P_{\ell}(\cos\theta) \end{aligned}$$

$$A_{\ell}(r) = c_{\ell}^{(1)} h_{\ell}^{(1)}(r) + c_{\ell}^{(2)} h_{\ell}^{(2)}(r), \quad \text{where}$$

$$h_{\ell}^{(1)}(r) = j_{\ell} + i n_{\ell} \quad h_{\ell}^{(2)}(r) = j_{\ell} - i n_{\ell}$$

are Hankel functions of the first and second kind and j_{ℓ}, n_{ℓ} are spherical bessel, neumann functions.

A bit of math (which I will leave out for now) gives a general expression

$$\tan \delta_{\ell} = \frac{\alpha_{\ell} j_{\ell}'(ka) - \beta_{\ell} j_{\ell}(ka)}{\alpha_{\ell} n_{\ell}'(ka) - \beta_{\ell} n_{\ell}(ka)}, \quad \text{where}$$

β_{ℓ} is the logarithmic derivative of A_{ℓ} determined by continuity at the boundary of the scattering region. For the special case of interest here, the wave function must vanish at the boundary. So, we have

$$j_{\ell}(ka) \cos \delta_{\ell} - n_{\ell}(ka) \sin \delta_{\ell} = 0$$

$$\tan \delta_{\ell} = \frac{j_{\ell}'(ka)}{n_{\ell}'(ka)}$$

$$\tan \delta_0 = \frac{j_0'(ka)}{n_0'(ka)} = \frac{\sin ka/ka}{-\cos ka/ka} = -\tan ka$$

$$\therefore \delta_0 = -ka$$